## Combinatorial and Computational Aspects of the Monomer-Dimer Problem: Extended abstract

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The exponential growth rate h (with respect to the natural logarithm) of the number of configurations on a multi-dimensional grid arises in the theory of various phenomena [3]. In physics  $e^h$  is viewed as the entropy (per atom) of the corresponding "hard model"; in mathematics h is called the topological entropy [5]; and in information theory h (with respect to  $\log_2$ ) is called the multi-dimensional capacity [12]. In the 1-dimensional case  $e^h$  is equal to the spectral radius  $\rho(\Gamma)$  of certain digraph  $\Gamma$ . There are very few 2-dimensional model where the value of h is known in closed form. In [7] we give a complete up-to-date theory of the computation of h by using lower and upper bounds that converge to h. It refines the techniques described in [6] by using an automorphism subgroup of a certain graph. As a demonstration of these techniques, we compute the topological entropy of the monomer-dimer covers of the 2-dimensional grid to 8 decimal digits and of the 3-dimensional grid with an error smaller than 1.35%.

Let  $\mathbb{Z}^d$  be the grid of integer points in d-dimensional space  $\mathbb{R}^d$ . A dimer is a domino consisting of two neighboring atoms occupying the places  $\mathbf{i}, \mathbf{i} + \mathbf{e}_k \in \mathbb{Z}^d$ . A monomer is a single atom occupying the place  $\mathbf{i} \in \mathbb{Z}^d$ . A monomer-dimer cover, respectively dimer cover, of  $\mathbb{Z}^d$  is a partition of  $\mathbb{Z}^d$  into monomers and dimers, respectively dimers. We denote by  $h_d$  and  $h_d$  the entropies of the monomer-dimer and dimer covers, respectively. In other words,  $h_d$  is the limit of the logarithm of the number of monomer-dimer covers of a box in  $\mathbb{Z}^d$  divided by the volume of the box, as the dimensions of the box grow to infinity; and similarly for  $h_d$ . It is straightforward to compute the values  $h_1 = \log \frac{1+\sqrt{5}}{2}$  and  $h_d = 0$ . The big breakthrough in the sixties was a close formula for  $h_d$  in [2]. The exact values of  $h_d$  for  $d \geq 2$  and  $h_d$  for  $d \geq 3$  are unknown.

It was shown in [8] that for  $p \in [0,1]$ , there exists the entropy  $\lambda_d(p)$  of the monomerdimer covers of  $\mathbb{Z}^d$ , where p is the "density" of dimers, i.e., the number of dimers in the cover divided by one half of the volume. The entropy  $\lambda_d(p)$  is a continuous concave function of p and  $\lambda_d(1) = \widetilde{h}_d$ . We show that  $h_d = \max_{p \in [0,1]} \lambda_d(p)$ . The van der Waerden conjecture for permanents of doubly-stochastic matrices gives a lower bound on  $\widetilde{h}_d$ . The improved lower bound for the permanents of 0-1 matrices [11] gives the currently best lower bound  $\widetilde{h}_3 \geq 0.440075$ . A recent breakthrough [1] gives the upper bound  $0.463107 \geq \widetilde{h}_3$ , improved to  $0.457547 \geq \widetilde{h}_3$  by Lundow [10].

In [7] it is to shown that the entropies  $h_d$  and  $h_d$  obey upper and lower bounds similar to the upper and the lower for the entropy of configurations with the symmetric isotropic

nearest neighbor graph. The bounds for  $h_d$  are stated in terms of the spectral radii of certain multigraphs whose automorphism group has a subgroup isomorphic to the the group of rigid motions of the (d-1)-dimensional torus  $(\mathbb{Z}/m_1\mathbb{Z})\times\cdots\times(\mathbb{Z}/m_{d-1}\mathbb{Z})$ . This fact enables us to compute the values of  $h_2$  and  $h_3$  with good precision. We also show that  $\lambda_d(p)$  can be bounded below by using the generalized van der Waerden conjecture (Tverberg's conjecture), proved by the first author in [4]. In particular, these lower bounds yield a lower bound for  $h_d$ . For d=2 this lower bound is somewhat weaker than the one obtained from the numerical computations, but for d=3 the situation is reversed.

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