

Folding during soft-sediment deformation



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Abstract: The detailed analysis of folding in rocks was in part pioneered by John Ramsay, and resulted in a range of techniques and criteria to define folds. Although folding of unlithified or 'soft' sediments is typically assumed to produce similar geometries to those in 'hard rocks', there has to date been little detailed analysis of such folds. The aim of this paper is therefore to investigate folds developed during soft-sediment deformation (SSD) by applying techniques established for the analysis of tectonic folds during hard-rock deformation (HRD). We use the Late Pleistocene Lisan Formation exposed around the Dead Sea as our case study, as the laminated lake sediments record intricacies of fold detail generated during seismically triggered slumping of mass transport deposits (MTDs) towards the depocentre of the basin. While it is frequently assumed that folds created during SSD are chaotic and form disharmonic structures, we provide analyses that show harmonic fold trains may form during slumping, although larger upright folds cannot be traced for significant distances and are more typically disharmonic. Our analysis also reveals a range of fold styles, with more competent detrital-rich layers displaying buckles (Class 1B), as well as upright Class 1A folds marked by thickened limbs. Class 1A buckle folds are generally considered to be created by flattening that overprints folds with an original Class 1B geometry. As thickened fold limbs are truncated by overlying erosive surfaces, the vertical flattening is considered to have occurred during the slump event. Different fold shapes may partially reflect variable flattening, depending on the original orientation of upright or recumbent folds, together with continued downslope-directed simple-shear deformation that modifies the fold geometry. Analysis of fold wavelength, amplitude and bed thickness allows us to plot strain contour maps, and indicates that beds defining slump folds display viscosity contrasts in the range of 50–250, which are similar to values estimated from folds created during HRD in metamorphic rocks. A range of re-fold patterns, similar to those established by John Ramsay in metamorphic rocks, are observed within slumps, and are truncated by the overlying sediments, indicating that they formed during a single progressive slump event rather than distinct 'episodes' of superimposed deformation. This study confirms that techniques developed for the analysis of folds created during HRD are equally applicable to those formed during SSD, and that resulting folds are generally indistinguishable from one another. Extreme caution should therefore be exercised when interpreting the origin of folds in the rock record where the palaeogeographical and tectonic contexts become increasingly uncertain, thereby leading to potential misidentification of folds created during SSD.

It is generally assumed that slump folds generated during soft-sediment deformation (SSD) within mass transport deposits (MTDs) bear a close resemblance to tectonic folds created during hard-rock deformation (HRD) (see discussions in Elliot & Williams 1988; Maltman 1994; Alsop *et al.* in press a). However, there has in reality been very little systematic analysis of fold styles created during SSD, with some notable exceptions including Woodcock (1976a, b, 1979), Farrell & Eaton (1987), Waldron & Gagnon (2011) and Ortner & Kilian (2016).

Hudleston (1986, p. 238) noted that 'shape alone cannot be used to distinguish soft-sediment folds from tectonic folds' (see also Woodcock 1976b). Comparison of curvilinear fold patterns in metamorphic rocks and soft sediments also reveals no significant geometric difference between aspects of curvilinear fold geometries (e.g. Alsop *et al.* 2007).

While folding within metamorphic rocks is generally created by deformation associated with recrystallization, folding within unlithified or 'soft' sediments is typically achieved via independent

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particulate flow, where individual grains move relative to one another to create a range of structures (e.g. Knipe 1986). The exact nature of these structures is dependent on the ratio of pore fluid pressure to cohesive strength of the sediment (e.g. Knipe 1986; Ortner 2007); with folds and shears being created where fluid pressure is less than cohesive strength, leading to hydroplastic deformation. The examination of fold geometries is important as Waldron & Gagnon (2011) realized that fold styles in different sand and mud layers may be used to distinguish the degree of lithification during the fold process, and thereby identify ‘soft-sediment’ folds. Despite the obvious differences in the mechanisms of folding within HRD and SSD noted above, the geometry of the resulting structures are so similar to one another that criteria to distinguish between these different types of folds in ancient settings should be considered (see Alsop *et al.* *in press a* for a review).

The broad aim of this contribution is to document and investigate fold styles created during SSD by applying classical techniques of structural geology such as dip-isogon analysis and examination of fold-interference patterns established by John Ramsay, among others (e.g. Ramsay 1967; Ramsay & Huber 1987). Van der Pluijm & Marshak (2004, p. 25) discussed slump folds and noted that ‘folds in one layer are of a different size and orientation than the structures in adjacent layers’, suggesting a largely disharmonic style. We therefore examine trains of soft-sediment folds to determine whether they are, indeed, disharmonic; using techniques of fold spacing compared to distance they can be traced along their axial surface (e.g. Twiss & Moores 2007, p. 290). We also analyse folded layer thickness, amplitude and wavelength in an attempt to estimate the viscosity contrast between layers during sediment folding together with the percentage of shortening (e.g. Schmalholz & Podladchikov 2001). Our analysis aims to help answer the following research questions:

- Does slumping create harmonic or disharmonic folds?
- Are re-fold patterns in slumps similar to those in metamorphic rocks?
- What range of fold styles are created during slumping?
- How are slump folds subsequently modified?
- Are estimates of viscosity contrasts in slump folds similar to those in metamorphic rocks?

We first describe the regional setting of the case study area in the Dead Sea Basin, before analysing folds created during SSD linked to downslope-directed slumping of MTDs towards the basin.

Regional setting

The Dead Sea Basin is a pull-apart structure on the Dead Sea Fault (DSF), which is marked by two major parallel fault strands that generate numerous earthquakes (Fig. 1a, b) (e.g. Marco *et al.* 1996, 2003; Ken-Tor *et al.* 2001; Migowski *et al.* 2004; Begin *et al.* 2005). This transform is thought to have been active from the Miocene to recent, including during the deposition of the Late Pleistocene (70–15 ka) Lisan Formation that forms the focus of the present study (e.g. Bartov *et al.* 1980; Garfunkel 1981; Haase-Schramm *et al.* 2004). The Lisan Formation comprises millimetre-scale aragonite laminae that were precipitated from the upper waters of Lake Lisan during the summer, together with more detrital-rich layers washed into the lake during flood events (Begin *et al.* 1974) (Fig. 2a, b). The detrital units were sampled <1 km NE of the Peratzim site by Haliva-Cohen *et al.* (2012) (see Fig. 1b), and compositionally consist of quartz and calcite grains with minor feldspar, and clays (illite–smectite). Detrital laminae within the varved aragonite-rich Lisan Formation display grain sizes of *c.* 8–10 μm (silt), while the thicker (>10 cm) detrital-rich units are generally coarser grained (60–70 μm) and can be classified as very fine sands (Haliva-Cohen *et al.* 2012). Although deposited on slopes of <1°, the Lisan Formation contains numerous intraformational fold and thrust horizons that are capped by undeformed beds and are considered to be seismically-triggered (Alsop & Marco 2013; Alsop *et al.* 2016; Lu *et al.* 2017) (Fig. 2a). Some evidence for bedding-parallel shear during seismic events is also preserved by offset clastic dykes (Weinberger *et al.* 2016).

The Lisan Formation is exposed for *c.* 100 km along the western margin of the Dead Sea Basin, and displays a systematic variation in the orientation of slump fold and thrust systems within it. In the northern portions slumping is directed towards the ESE, in the central portion around Masada slumping is towards the east, whilst in the southern area around Peratzim slumping is NE directed (Alsop & Marco 2012a) (Fig. 1b). Combined with westerly-directed slump folds recorded from the eastern shore of the Dead Sea in Jordan (El-Isa & Mustafa 1986), this suggests a regional pattern of radial slumping towards the depocentre of the Dead Sea Basin (Alsop & Marco 2012a). In the extreme southern area around Zin, slumping is towards the south and is interpreted to be influenced by the nearby NW–SE-trending Amazyahu Fault (Weinberger *et al.* 2017; Alsop *et al.* 2018, *in press b*) (Fig. 1b). This overall pattern has been subsequently corroborated by anisotropy of magnetic susceptibility (AMS) fabrics that support a radial pattern of MTD slumping (Weinberger *et al.* 2017; Levi *et al.* 2018).

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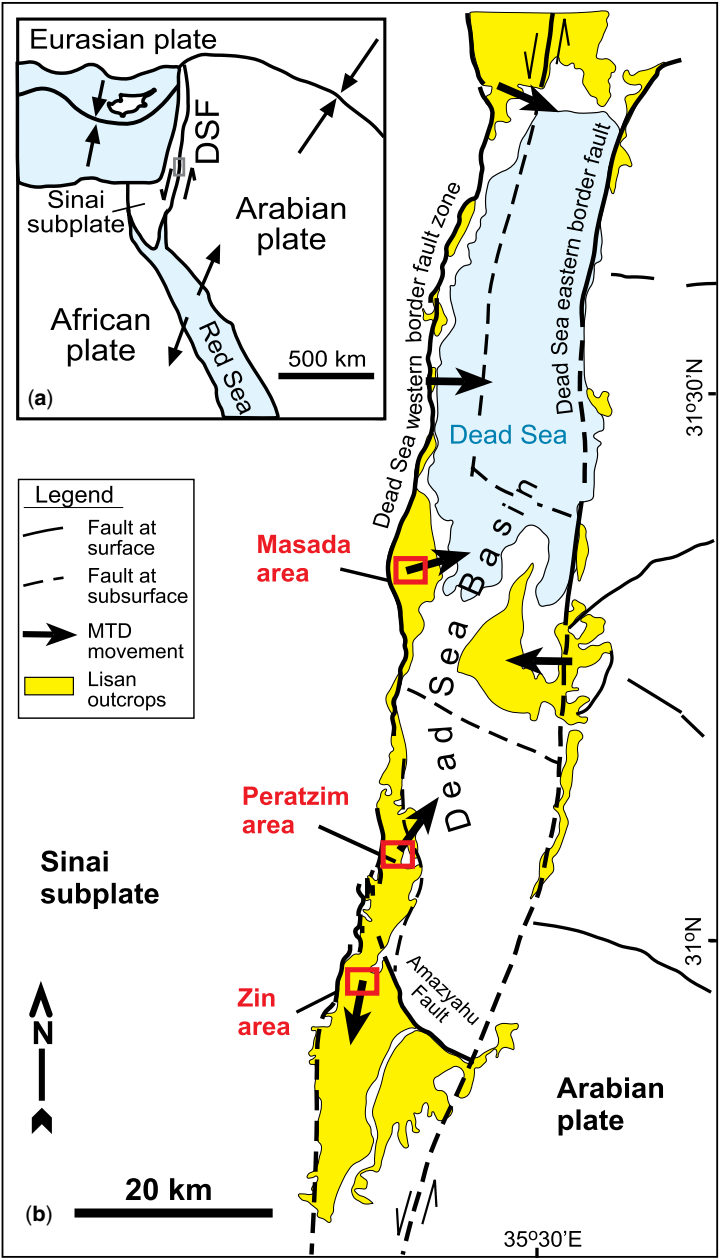
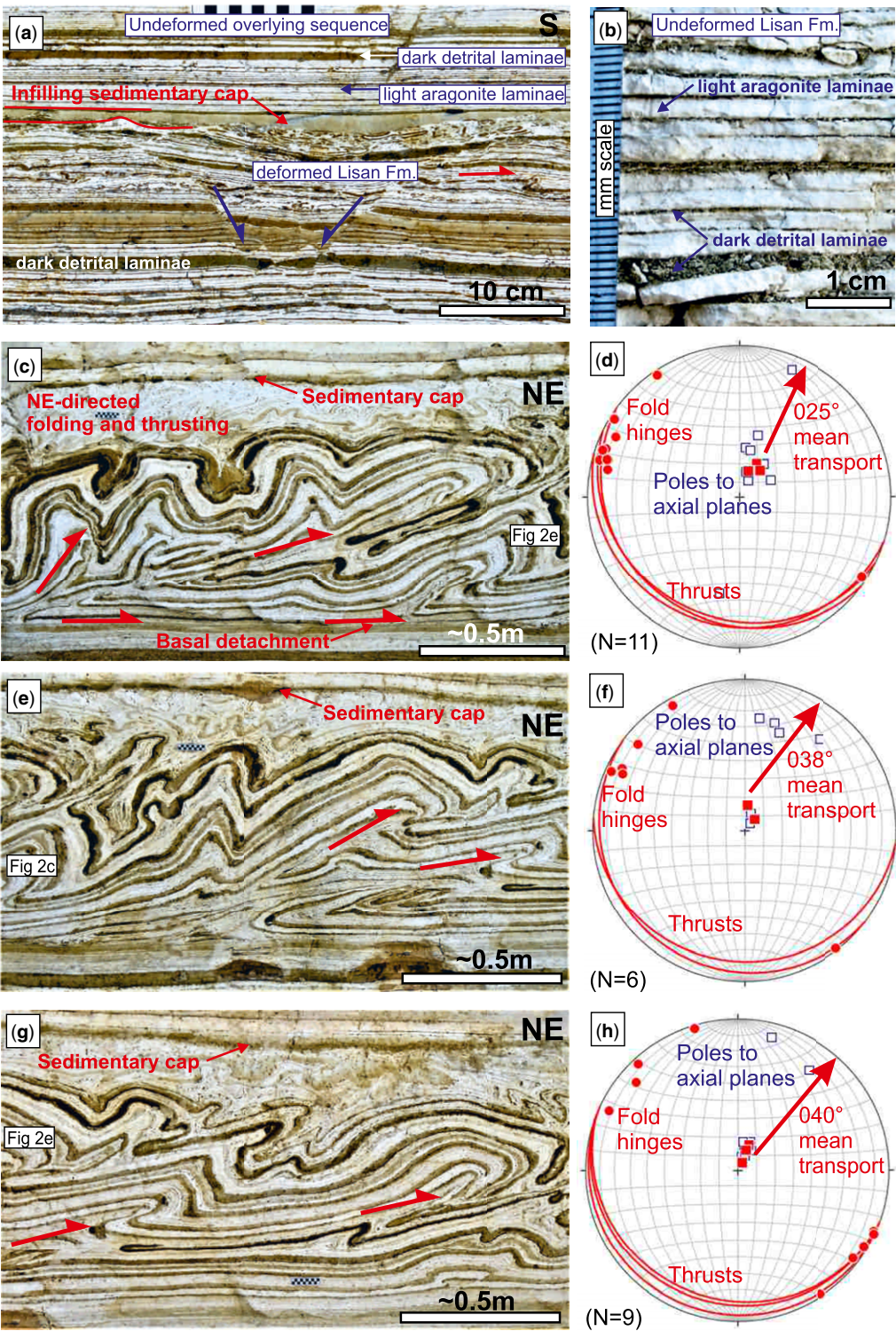


Fig. 1. (a) Tectonic plates in the Middle East. General tectonic map showing the location of the present Dead Sea Fault (DSF). The Dead Sea Fault transfers the opening motion in the Red Sea to the Taurus–Zagros collision zone. (b) Generalized map (based on Sneh & Weinberger 2014) showing the current Dead Sea, including the position of the Masada, Peratzim and Zin localities referred to in the text. The extent of the Lisan Formation outcrops are also shown, together with the general slump directions of the MTDs around the basin.

The Dead Sea Basin is an ideal place to study structures associated with SSD as the relatively cohesive muds and precipitated aragonitic layers define an intricate varve-like stratigraphy that define

a range of detailed structures that may not survive elsewhere (Fig. 2a, b). The MTD horizons and reworked zones of sediment we are particularly interested in for this case study are best exposed



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around Masada (31° 1847.0' N, 35° 2228.3' E), Peratzim (31° 0449.6' N, 35° 2104.2' E) and Zin (31.00615° N, 35.26342° E) in the central and southern Dead Sea areas, respectively (Fig. 1b) (see Weinberger *et al.* 2017). All of these sites are positioned c. 1 km east of the Dead Sea western border fault zone, with Cenomanian–Senonian carbonates preserved further to the west in the footwall to this fault. For most of the time between 70 and 28 ka, Lake Lisan in these areas had a maximum depth of 100 m or less, apart from a brief period from 26 to 24 ka when water was up to 200 m deep (Bartov *et al.* 2003).

Analysis of slump fold and thrust displacement

Six individual slump sheets that form MTDs have been recognized in the Peratzim area (Alsop *et al.* 2016). Although some variability exists, slump fold hinges are typically NW–SE trending, and verge and face towards the NE, with transport directed towards the depocentre of the basin (Alsop *et al.* 2016) (Fig. 2c–h). Fold axial planes dip variably towards the SW, while downslope-verging fore thrusts dip towards the SW (Fig. 2c–h) (Alsop *et al.* 2017a). At the Masada locality, fold hinges are typically NNW–SSE trending with gently WSW-dipping axial planes and overall ENE-directed vergence, with transport towards the depocentre of the basin (Alsop & Marco 2012a) (Fig. 3a, b). At the Zin locality, fold hinges are east–west trending, and typically verge and face towards the south (Weinberger *et al.* 2017; Alsop *et al.* 2018). As with previous studies, care has been taken that cuttings along wadi walls are developed at high angles to fold hinges, thereby providing transport-parallel (or hinge-normal) profile sections to undertake structural analysis (e.g. Alsop *et al.* 2017b).

Each MTD horizon is overlain by undeformed sediments that were deposited after each slope failure event (e.g. Fig. 2a). The tops of MTDs, and the folds they contain, may be truncated by the overlying sedimentary ‘cap’ that was deposited out of suspension following the failure event (Alsop & Marco 2012b; Alsop *et al.* 2016, *in press a, b*). This cap, which may be up to 20 cm thick, blankets and infills

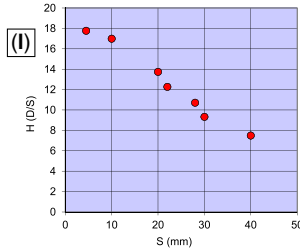
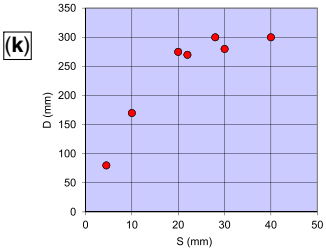
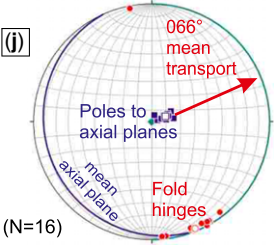
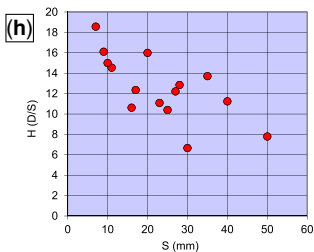
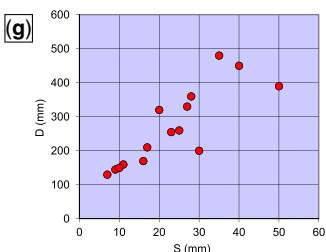
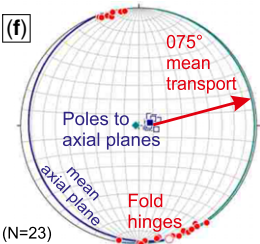
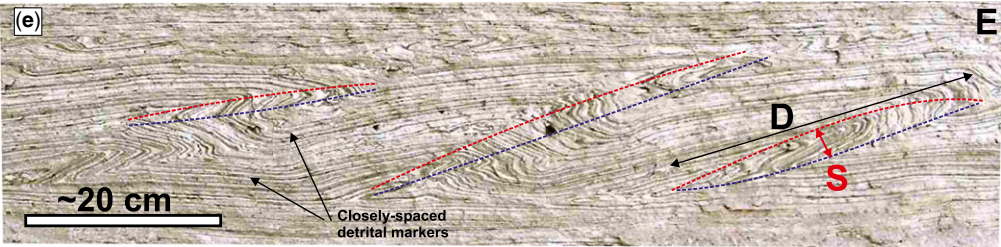
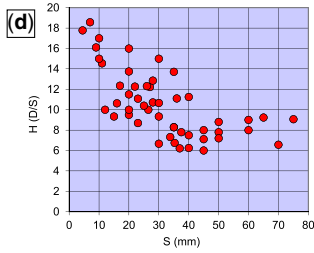
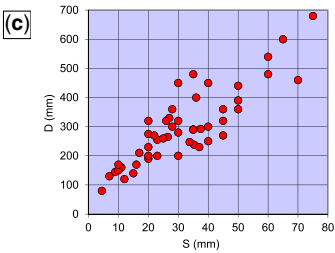
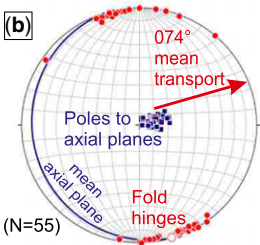
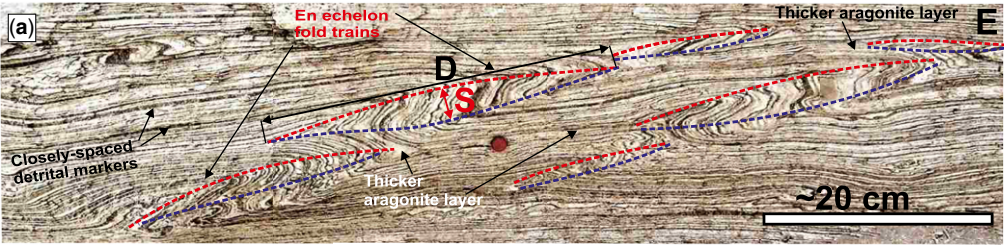
the underlying slump topography, and demonstrates that the structures formed at or close to the sediment surface. The studied sections form in the upper part of the Lisan Formation and have never had significant overburden (<10 m) above them. As such, the structures created during slumping have not undergone significant later modification, and preserve pristine geometries unmodified by later compaction (Alsop *et al.* 2017a, b, *in press a*). In addition, the aragonite- and detrital-rich alternations within the Lisan Formation form a bilaminate sequence: that is, ‘comprising only two different types of layers which alternate with each other’ (Price & Cosgrove 1990, p. 307). Such bilaminates simplify the analysis of folding, although they need not be regularly spaced (equal thickness) and may form either multilayer packages of folding (where folding is closely spaced) or single-layer buckles where competent beds are separated from one another by a weaker ‘host’.

Harmonic and disharmonic slump folding

A harmonic fold is simply defined as being continuous along its axial surface for ‘many multiples of the half wavelength’, whereas a disharmonic fold ‘dies out within a couple of half wavelengths’ (Twiss & Moores 2007, p. 289; see also Fossen 2016, p. 259). The half wavelength may be approximated by the spacing (S) of adjacent axial surfaces, while the continuity of the axial surface may be directly measured (D) to produce a ratio (H) which is equivalent to D/S (see Twiss & Moores 2007, p. 290) (Fig. 3a).

Upright folds within the Lisan Formation are frequently disharmonic, as they are positioned above basal detachments and are overlain by sedimentary caps meaning that they cannot be traced for any significant distance (D) along their axial surfaces (e.g. Fig. 2c, e, g). However, we have also analysed trains of recumbent slump folds (Fig. 3a, e, i). In general, the maximum distance (D) that individual decimetric-scale folds may be traced along their axial planes increases as the maximum spacing (S) between adjacent axial planes becomes greater (Fig. 3a, c). In addition, the ratio H ranges from c. 18 to c. 6, and is shown to reduce in a non-linear way as S increases (Fig. 3d). Analysis of individual fold

Fig. 2. (a) & (b) Photographs of light-coloured aragonite-rich laminae and dark detrital-rich laminae forming the Lisan Formation. An infilling sedimentary cap that blankets the underlying structures at Zin is shown in (a), while (b) is a close-up image of undeformed laminae from Peratzim. Refer to Figure 1b for details of the locations. Pairs of photographs (c), (e) & (g) and associated stereonet (d), (f) & (h) of NE-verging fold and thrusts developed within slump 4 at Peratzim (see Alsop *et al.* 2016). In stereonets, fold hinges (solid red circles), poles to axial planes (open blue squares), and associated thrusts are shown as red great circles and poles as solid red squares. Calculated slump transport is towards the NE (see Alsop & Marco 2012a). All photographs marginally overlap with one another and have NE on the right-hand side for consistency.



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systems reveals broadly similar patterns of D , S and resulting H (Fig. 3e–l). Individual fold packages may form en echelon fold trains that migrate up through the folded sequence where they cross thicker aragonite-rich units (Fig. 3a). These fold packages do not directly interact with one another and do not therefore create re-fold patterns (see the following section). Overall, these recumbent folds may be described as broadly harmonic, as $D \gg S$.

Refold patterns within individual slump horizons

Working on polyphase deformed metamorphic rocks in the Scottish Highlands, Ramsay (1962, 1967, p. 518) developed a classification scheme of refolded folds, based on the relative orientations of different sets of fold hinges and axial surfaces that were superimposed on one another, and considered to develop sequentially during orogenesis. Ramsay's (1962, 1967) classic Loch Monar dam outcrops comprise superb water-washed exposures in psammities of the Precambrian Moine Supergroup, and permit intricate details of folding to be examined (e.g. Fig. 4a). Ramsay (1967, p. 521) recognized Type 1 or 'dome and basin' refold patterns created by superimposition of fold hinges and axial planes at high angles (c. 90°) to one another (Fig. 4b, c). Type 2 crescent or 'angel-wing/mushroom' styles of refolding formed by superimposition of fold hinges and axial planes at moderate to high angles to one another were also observed (Ramsay 1967, p. 525); as were Classic Type 3 'hook' refolds that form where superimposed fold hinges are coaxial, while axial surfaces are at high angles to one another (Ramsay 1967, p. 530) (Fig. 4d). This general classification scheme has subsequently been applied across a huge range of metamorphic rocks, scales and settings.

Within the Peratzim area of the Dead Sea case study (Fig. 1b), Type 1 or 'dome and basin' refold patterns are developed in wadi sections within individual MTDs (Fig. 4e, f). Type 2 crescent or 'angel-wing/mushroom' styles of refolding are also observed (Fig. 4g), as are classic Type 3 'hook' refolds that are particularly common (Fig. 4h, i, j). Refold patterns

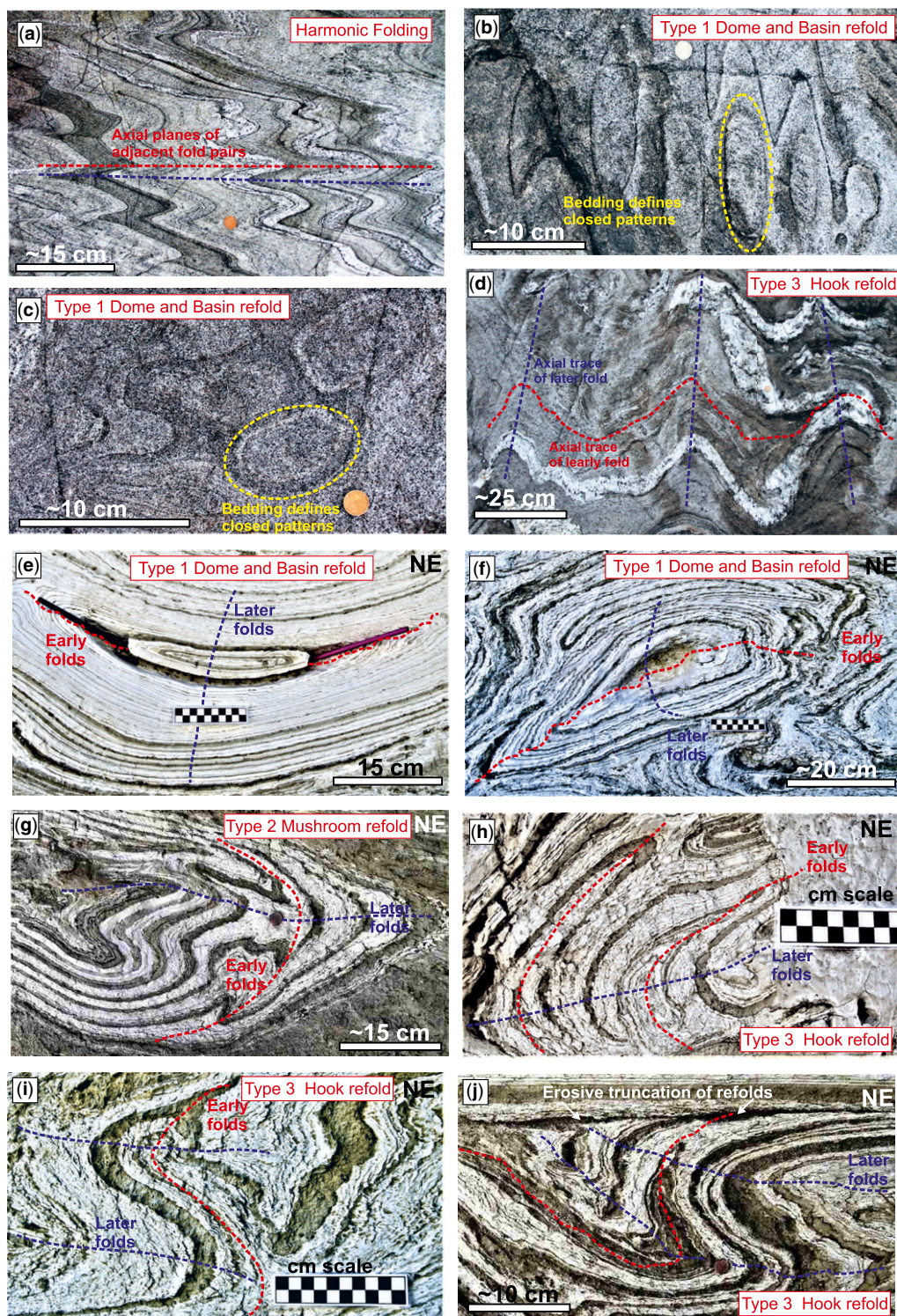
are truncated by the overlying sedimentary cap that was deposited from suspension following the slump event (Alsop & Marco 2012b) (e.g. Fig. 4j). This indicates that refolding was completed during the slump event and was not a much later re-working created via subsequent loading or later slumping. Thus, a re-working or modification of folds to create classic refold patterns occurs during a single, progressive slump event.

Dip-isogon analysis around slump folds

The dip-isogon method is a well-established technique of fold classification in lithified rocks developed by John Ramsay (e.g. Ramsay 1967, p. 363). In this method, dip isogons join points of equal dip on adjacent folded surfaces within the fold profile, t_0 is the layer thickness measured along the axial surface, while t_α is the orthogonal layer thickness measured at various angles (α) to the reference plane (Fig. 5a). Graphs normalize thicknesses by using t'_α (where $t'_\alpha = t_\alpha/t_0$) and plot this value against dip angle (α) to create a series of fold classes (Ramsay 1967, p. 366). Class 1 folds are marked by convergent dip isogons (Fig. 5b), Class 2 folds by parallel dip isogons (Fig. 5c) and Class 3 folds by diverging dip isogons (Fig. 5d) (e.g. Ramsay 1967, p. 365; see Fossen 2016, p. 263).

We use the dip-isogon method to analyse and compare fold geometries formed in aragonite-rich and detrital-rich units within the case study at Peratzim (Fig. 1b). The unlithified sediment enables easy excavation, and reveals that east–west-trending asymmetrical buckle folds with gently south-dipping axial planes verge towards the north (Fig. 5e, f, g). Our analysis includes data from both the upper and lower limbs of folds, and shows that folds within aragonite-rich units display gently convergent to parallel dip isogons that typically define Class 1C folds (Ramsay 1967, p. 367; Fossen 2016, p. 263) (Fig. 5h, i). Conversely, the detrital units are marked by a strongly convergent isogon pattern representing Class 1C or 1B parallel folds consistent with buckling (Fig. 5h, i). In detail, the upper limbs are in part Class 1C (Fig. 5h, i), while the lower overturned limb may exhibit a component of thickening

Fig. 3. Harmonic and disharmonic fold analysis from folds in Masada (see Fig. 1b for the location). (a) Recumbent en echelon folds verging towards the east. The half wavelength of folds is approximated by the spacing (S) of adjacent axial surfaces (shown in red and blue), while the continuity of the axial surface may be measured directly (D) to produce a ratio (H) that is equivalent to D/S (see Twiss & Moores 2007, p. 290). (b) Stereonet of folds from Masada shown in (a) with the transport direction towards 074°; the key to the stereonet is given below. (c) Graph showing D plotted against S . (d) Graph showing H plotted against S . (e) Photograph of an individual Masada locality with stereonet (f), and plots of H , D and S parameters (g) & (h). (i) Photograph of an individual Masada locality with stereonet (j), and plots of parameters (k) & (l). In all stereonets, fold hinges (solid red circles), mean fold hinge (open red circle), poles to axial planes (solid blue squares), mean axial plane (blue great circle and pole as an open blue circle). The great circle and pole to subhorizontal bedding are shown in green in (f) & (j).



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relative to the hinge area to create Class 1A folds (Figs 5h, i & 6a, b). However, as folds become progressively more asymmetrical and overturned, the lower limb is consequently thinned (Fig. 5e). These overturned buckles also display a reduction in the angle of axial-planar dip from 70° to 24°, together with a slight re-orientation of fold hinges from 092° towards 063° and the general slope direction at Peratzim (Figs 2c–h & 5e, f, g). In general, more upright folds within detrital-rich units display Class 1A or Class 1B (parallel) buckle fold styles, while recumbent folds are marked by Class 1C or Class 2 (similar) geometries (Fig. 6a–h). Aragonite-rich units generally display a parallel isogon pattern most consistent with Class 2 similar folding (Fig. 6c, d, g, h). These relationships collectively indicate that detrital layers were locally more competent (see also Alsop *et al.* 2017a, b), and suggest that folding involved a component of progressive deformation that resulted in tightening and rotation of fold hinges towards the flow direction, while axial planes were rotated into the plane of flow.

Ramsay (1967, p. 432) recognized that multilayer folds may be composed of adjacent beds defining different classes of folds, such as Class 1C in more competent units and Class 3 in the incompetent horizons, although the overall fold represents the sum of these layers and will closely resemble a Class 2 similar fold. Within individual folds of the Lisan Formation, different layers define different classes of fold, with aragonite-rich layers displaying Class 2 (similar) folds, while adjacent detrital-rich beds are marked by Class 1C fold geometries (Fig. 6g, h). Such varying styles of folding allow an individual fold to propagate further along its axial plane without encountering significant accommodation problems around the fold hinge (e.g. Price & Cosgrove 1990, p. 320).

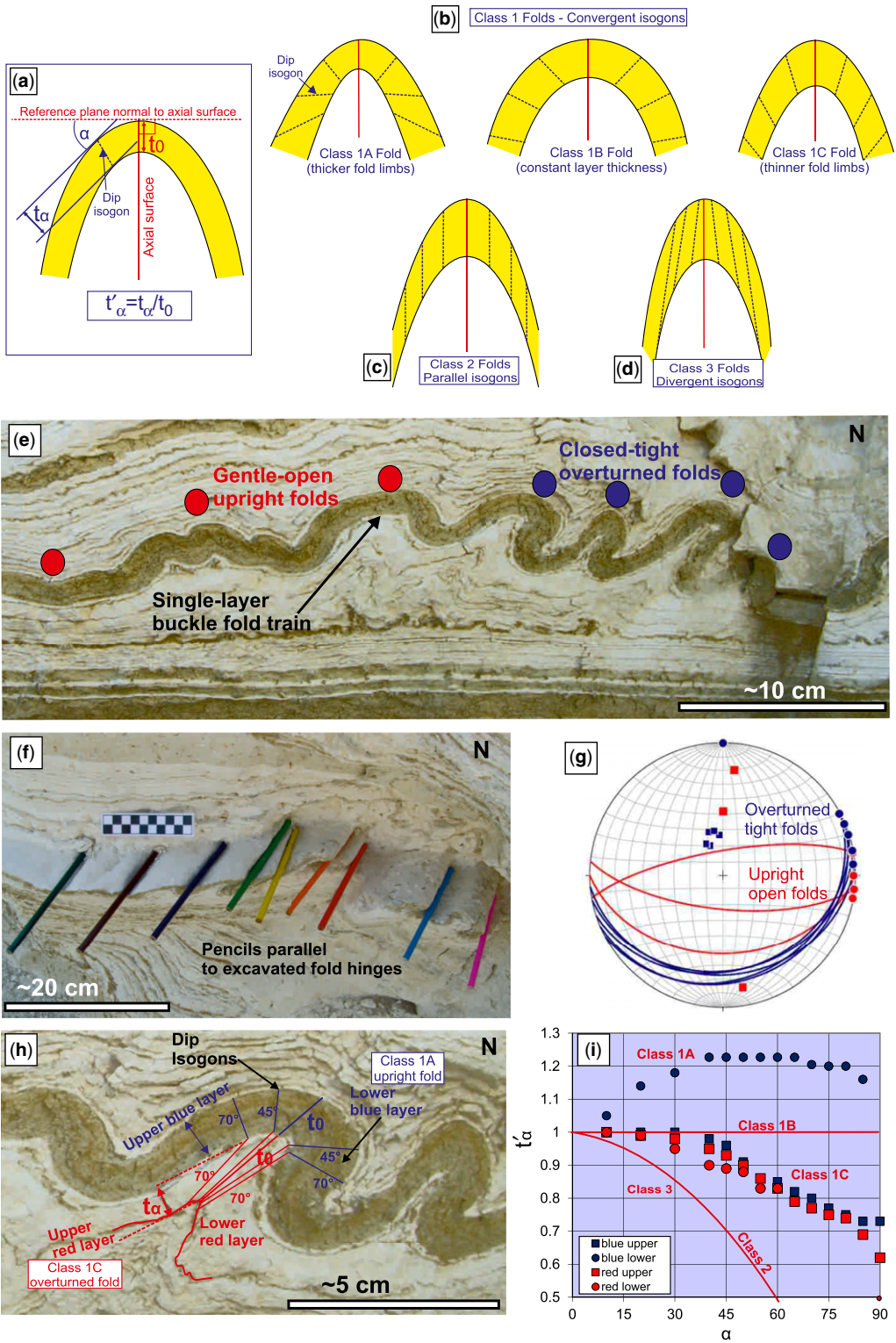
Post-buckle modification of slump folds

Buckling is simply defined as ‘the flexing or folding of a surface or series of parallel surfaces by a compressive stress directed along that surface or layer’ (Price & Cosgrove 1990, p. 273). However, Ramsay

(1967, p. 434) also recognized that buckle folds may undergo subsequent deformation and flattening that transforms their geometry from typical Class 1B parallel folds to Class 1C or Class 2 (similar) folds. A simple technique to determine the amount of post-buckle flattening was developed by Lisle (1992), and uses the inverse thickness method, where t is the orthogonal thickness measured between two tangents for the inner and outer layer boundaries. The inverse layer thickness ($1/t$) is then plotted for various orientations of the layer tangent around the fold (see Lisle 1992, p. 370) (Fig. 7a, b). This method has the significant advantage that post-buckle flattening does not have to be aligned parallel to the axial trace of the buckle fold (Lisle 1992). Although this technique was devised and applied to folds formed during HRD (e.g. Lisle 1992; Alsop *et al.* 1998, fig. 7), the present study forms its first use on folds created during SSD. Within the study area, elliptical ratios using the method of Lisle (1992) range between 1.77 (Fig. 7a, b), 1.70 (Fig. 7e, f) and 1.97 (Fig. 7i, j), suggesting a relative vertical shortening and flattening of c. 30%.

Within the Lisan Formation, folding of detrital-rich marker layers that are more competent than the surrounding aragonite-rich horizons (e.g. Alsop *et al.* 2017a, b) results in a parallel (Class 1B) style of folding (Fig. 5h, i). The flattening of upright Class 1B folds by vertical shortening results in thicker limbs compared to hinges (Class 1A) (Figs 6a, b & 7a–l), whereas the potential vertical flattening of recumbent horizontal folds results in limbs becoming thinner (relative to hinges) and the creation of Class 1C or similar (Class 2) folds (Fig. 6c–h) (see Farrell & Eaton 1987, fig. 11). Examples of both types of folds defined by detrital-rich marker horizons that are considered to have originally formed Class 1B folds are observed in the Lisan Formation (Fig. 6c–h). Clearly, the post-buckle flattening component is independent of the orientation of the buckle fold, with the orientation of the strain ellipse indicating that flattening was consistently subvertical (Fig. 7a–l). The thickening and thinning patterns observed around the slump folds are similar to the theoretical models of flattened folds as discussed by Twiss & Moores (2007, p. 374).

Fig. 4. Photographs (a–d) of minor folds within Moine psammities from the classic locality by Loch Monar Dam in NW Scotland [UK Grid NH1989 3882] used by Ramsay (1962, 1967) to establish a refold classification scheme. (a) Plan view of intense minor folding within interlayered pelites and psammities. Adjacent axial traces (highlighted in blue and red) define harmonic folds. (b) & (c) Plan view of Type 1 dome and basin fold interference patterns with bedding in psammite (highlighted in yellow) defining closed outcrop patterns. (d) Plan view of Type 3 hook fold interference patterns with axial traces of older folds (highlighted in red) being refolded by younger folds (highlighted in blue). Photographs (e–j) of sections through refolds within slumps at Peratzim (see Fig. 1b for the location). (e) & (f) Type 1 dome and basin refold pattern, with coloured pencils parallel to the fold hinges in (e). (g) Type 2 ‘mushroom’ refold pattern. (h), (i) & (j) Type 3 hook interference patterns. In (j), the refold is truncated by an overlying erosive surface and sedimentary cap. The 10 cm-long chequered rule and the 15 mm-diameter coin are for scale.



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An elegant development to the model of Lisle (1992) was provided by Srivastava & Shah (2006), who realized that strain ellipses (and the associated photograph of the fold) may simply be 'unstrained' on drafting packages by restoring the calculated strain ellipse back to an original circle: for example, Figure 7a and b are transformed to Figure 7c and d, respectively, when 'unstrained' (Fig. 7a–l). The upslope and downslope limbs have slightly different ellipse shapes, and this perhaps reflects the original limb thicknesses that vary as a result of fold vergence downslope towards the depocentre. The observation that upright Class 1A folds with thickened fold limbs are developed close to the sediment surface and are truncated by the overlying sedimentary cap (Fig. 7i–l) demonstrates that flattening formed during the actual slump event, and is not a later effect created by loading from overlying MTDs.

Analysis of layer thickness, amplitude and wavelength in slump folds

Layer thickness (h), amplitude (A) and wavelength (λ) of single-layer folds (rather than multilayers) may be measured in an attempt to estimate the strain and viscosity contrast between layers during folding (e.g. Schmalholz & Podladchikov 2001; see also Hudleston & Treagus 2010). The technique involves an analysis of single-layer folds (i.e. unaffected by neighbouring competent beds) and is based on results for linear viscous folding rather than power-law viscous folding (Schmalholz & Podladchikov 2001). The thickness of a layer (h) is measured orthogonal to the folded layer, while amplitude (A) is defined as half the distance from the trough to the crest of upright folds (e.g. Fig. 8a). Wavelength (λ) is defined as the distance between two points that occupy a similar position on the fold train (i.e. between adjacent synform hinges) (e.g. Fig. 8a). Schmalholz & Podladchikov (2001, p. 206) stated that wavelength may also be measured as double the horizontal distance between neighbouring fold

hinges (i.e. double the distance between antiform and synform fold hinges forming a fold pair). In this method of fold analysis, amplitude/wavelength (A/λ) is compared with layer thickness/wavelength (h/λ) on a strain contour map (Schmalholz & Podladchikov 2001). For any fold where amplitude, thickness and wavelength can be measured on the profile plane, estimates of bulk strain (in terms of % shortening) and the layer/matrix viscosity ratio can be made by reading the position of data directly off the map (e.g. Fig. 8a–r). This study forms the first use of this methodology in folds created during SSD, and assumes that all the layer shortening is taken up by buckling with no out-of-plane movement.

Within individual fold trains, the wavelength of folds (λ) reduces as amplitude (A) increases, so that the A/λ ratio defines a general trend when plotted against λ (Fig. 8a, b, d, e, g, h, j, k, m, n, p, q). In detail, when compared to wavelength, the A/λ ratio is not a straight line, with amplitude increasing more slowly than wavelength (Fig. 8e). In general, steeper folds typically have lower A/λ ratios, as recumbent folds can develop proportionally greater amplitudes due to being unconstrained and unhindered by the overlying sediment surface (e.g. compare upright folds in Fig. 8g & h with recumbent folds in Fig. 8p & q). The higher A/λ ratios of recumbent folds may also be a product of increased simple shear, which rotates the axial plane towards the shear plane (e.g. Figs 5e, f, g & 8a–c). Overall analysis of individual fold datasets on the strain contour maps (Schmalholz & Podladchikov 2001) reveals that most folded layers display viscosity contrasts in a range between 50 and 250, while calculated layer shortening is generally between 30 and 70% (Fig. 8c, f, i, l, o, r). Within individual fold trains, the percentage of contraction typically increases as folds become more inclined, or recumbent (e.g. Fig. 8a–c), with some small-scale recumbent folds recording values in excess of 70% (e.g. Fig. 8p–r). When plotted on strain contour maps, the overall trends of data from individual fold trains are slightly oblique to the established lines marking fixed

Fig. 5. (a) Dip isogons join points of equal dip on adjacent folded surfaces; t_0 is the layer thickness measured along the axial surface, while t_α is the orthogonal layer thickness measured at various angles (α) to the reference plane. (b) Class 1 folds are marked by convergent dip isogons, (c) Class 2 folds by parallel dip isogons and (d) Class 3 folds by diverging dip isogons. (e) Single-layer train of buckle folds defined by a detrital-rich bed at Peratzim (see Fig. 1b for the location). Buckle folds display a progressive tightening coupled with a rotation and reduction in dip of the axial surfaces towards the downslope direction. (f) Excavation of buckle folds shown in (e) reveals a 3D variation in fold-hinge orientations marked by different coloured pencils. The 10 cm-long chequered rule acts as scale. (g) Stereonet of upright folds (in red) and more overturned folds with inverted lower limbs (in blue) shown in (e) and (f). Axial planes are shown as great circles, and poles as solid squares. Fold hinges are shown by solid circles. (h) Close-up photograph of buckle folds that display thickened overturned limbs relative to the upper limbs (see e for the position). Representative 70° and 45° dip isogons are drawn on the upper and lower limbs of the fold. (i) t'_α graph (where $t'_\alpha = t_\alpha/t_0$) plotted against dip angle (α) to create a series of fold classes (Ramsay 1967, p. 366). The t'_α graph relates to the dip-isogon analysis of detrital-rich (blue) and aragonite-rich (red) layers around the buckle fold shown in (h). Data are divided into upper fold limbs (squares) and lower fold limbs (circles).

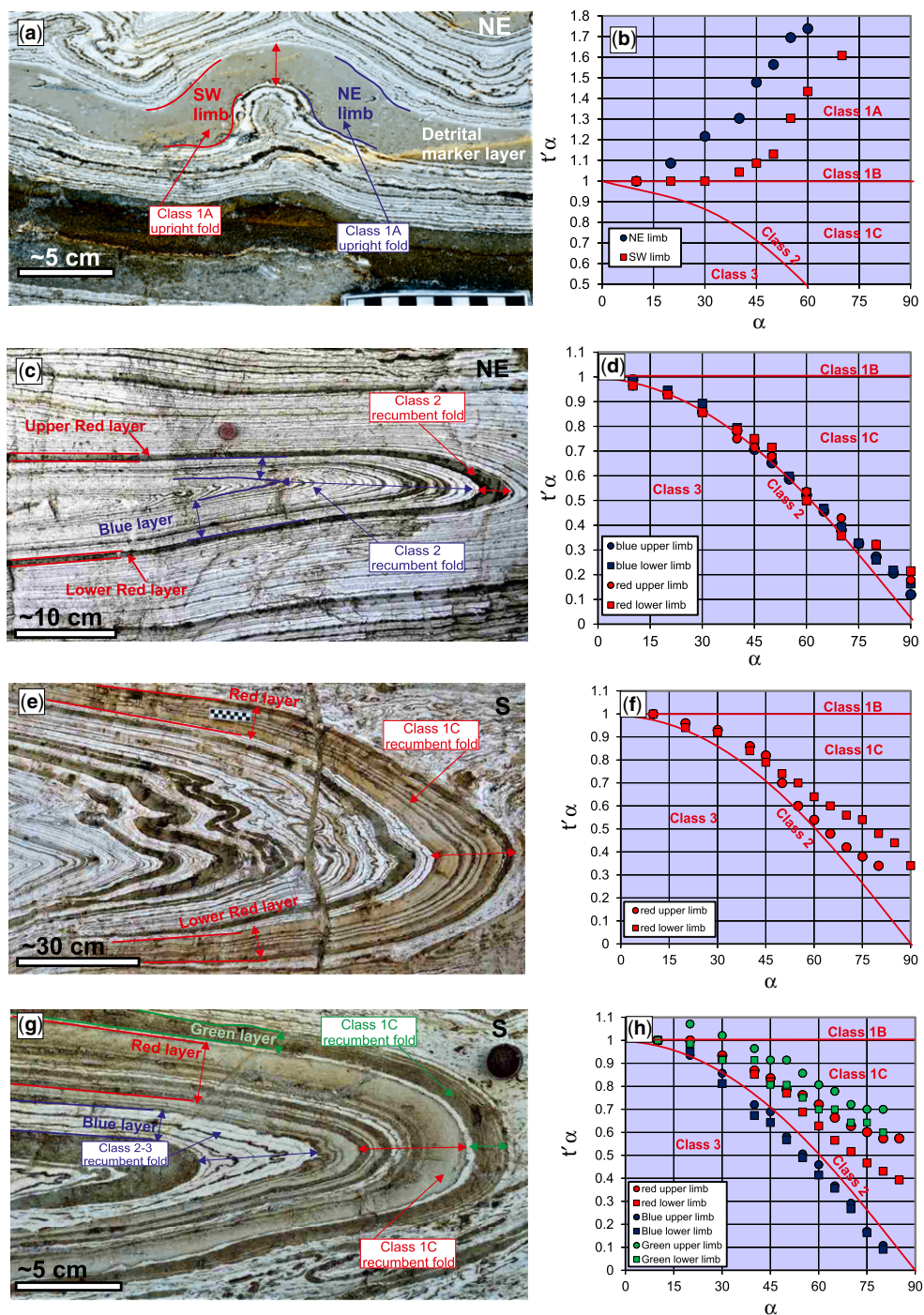


Fig. 6. Paired photographs and t'_α plots relating to the dip-isogon analysis of different layers around folds (a) & (b), (c) & (d), (e) & (f) and (g) & (h). Dip isogons join points of equal dip on adjacent folded surfaces; t_0 is the layer thickness measured along the axial surface, while t'_α is the orthogonal layer thickness measured at various angles (α) to the reference plane (see Fig. 5a). Upper fold limbs are represented by coloured circles and lower limbs by squares. Photographs (a) & (c) are from Peratzim, while (e) & (g) are from Zin.

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viscosity contrasts, with some plots suggesting that folds with a lower percentage of shortening are marked by lower viscosity contrasts compared to folds with a higher percentage of shortening (Fig. 8c). In other more typical cases, folds with a lower shortening percentage have greater viscosity contrasts compared to adjacent folds with a higher shortening percentage, resulting in more 'gentle' trends on strain contour maps (Fig. 8f, i, l, o).

When the combined dataset is considered, amplitude (A) shows a clear correlation with wavelength (λ) across a range of scales from millimetres (Fig. 9a), to centimetres (Fig. 9b) to larger folds shown on log-log plots (Fig. 9c). Similarly, layer thickness (h) and fold wavelength (λ) increase in tandem and correlate across a range of scales from millimetres to metres (Fig. 9d–f), as do layer thickness (h) and fold amplitude (A) (Fig. 9g, h). These general correlations produce a cluster of points on the overall strain contour plot, suggesting viscosity contrasts in the range of 50–250, and shortening of between 30 and 70% (Fig. 9i).

Discussion

Does slumping create harmonic or disharmonic folds?

Our results demonstrate that the distance (D) fold trains can be traced along their axial surfaces, compared to the spacing (S) with the neighbouring axial surface, shows distinct and consistent relationships (Fig. 3c, g, k). The ratio H (D/S) varies between c. 6 and 18, and generally increases as the spacing (S) reduces (Fig. 3d, h, i). These relationships collectively demonstrate that overturned and recumbent slump folds that can be traced for <1 m display harmonic fold relationships (Figs 3 & 10). Price & Cosgrove (1990, p. 307) suggested that the development of harmonic or disharmonic folds is dependent on the spacing of the competent beds, and suggested that 'adjacent layers must be closer than the sum of their dominant wavelengths (perhaps as little as 10% of this distance) before they will buckle harmonically'. Numerical modelling of linear viscous multilayer folding by Schmalholz & Mancktelow (2016, p. 1436) also indicates that more closely spaced layering may display 'contact strain' and will tend to fold harmonically. An inspection of Figure 3a reveals that the more competent detrital-rich markers are, indeed, typically closely spaced relative to the fold wavelength, and therefore supports this general model.

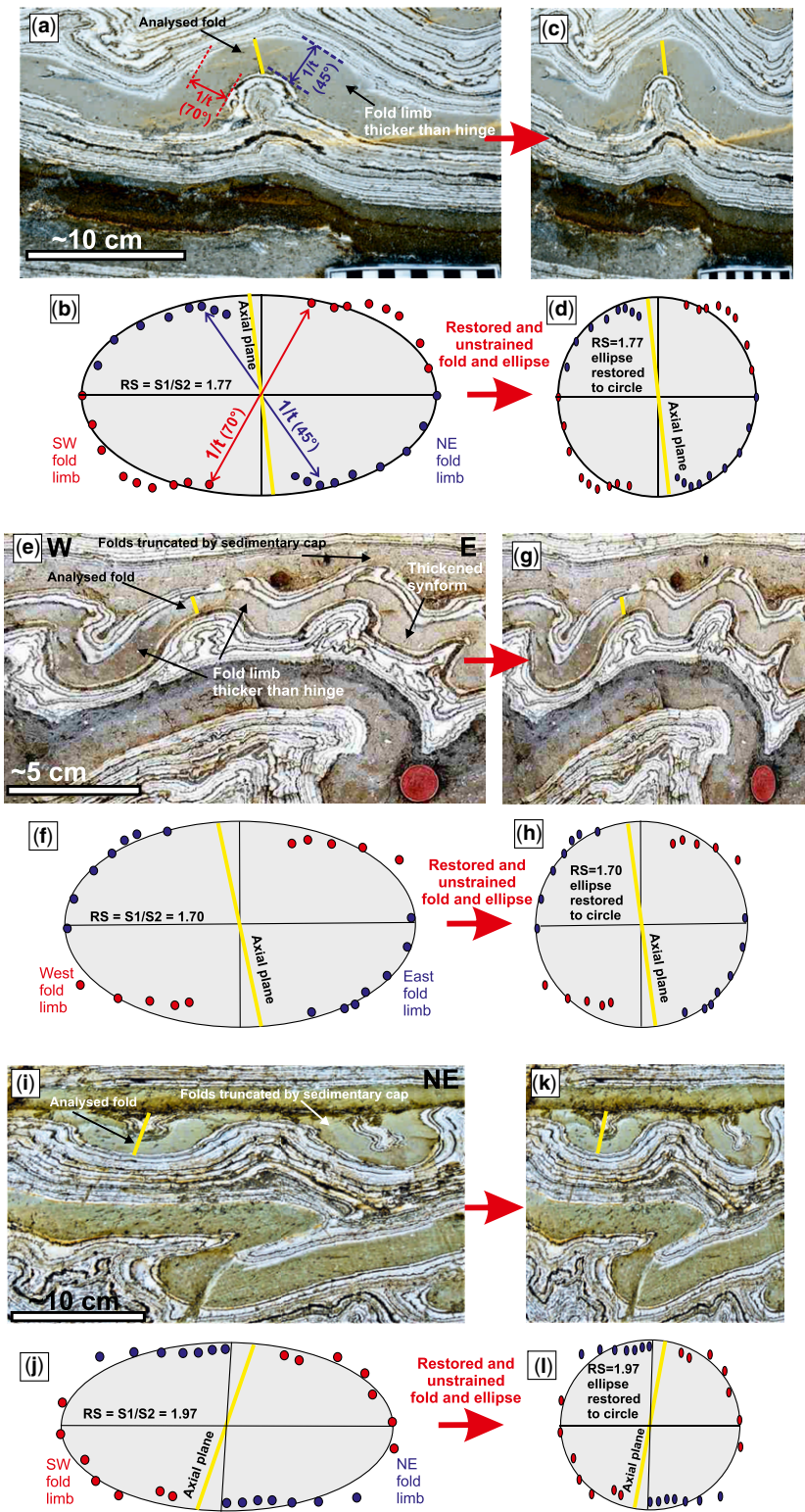
The observation that individual fold packages die out both up and down their axial surface (with associated reductions in the spacing of neighbouring axial surfaces) suggests that folds nucleated

in the central area of each fold package, and then propagated both upwards and downwards (Fig. 3a). Packages of en echelon folding generally transfer upwards through the slumped unit, with 'jumps' in each en echelon fold package typically occurring across thicker aragonite-rich units (Fig. 3a). Neighbouring en echelon systems tend to transfer across the same stratigraphic level marked by the weaker aragonite units (Fig. 3a), suggesting that they may form easy-slip horizons, and thereby control the position and geometry of fold trains. Although en echelon segments are broadly parallel, there is a general reduction in axial-surface dip together with interlimb angles up through the deformed sequence, such that the uppermost fold packages comprise subhorizontal tight isoclinal folds (Fig. 3a). This may suggest increasing shear towards the sediment surface during downslope-directed progressive deformation (Alsop & Marco 2013). The larger fold trains are preserved in the centre of the slump, perhaps indicating that deformation initiated here and then transferred downwards towards the underlying basal detachment. Such patterns are similar to those encountered within thrusts of the Lisan Formation, where larger displacement may occur on thrust ramps that initiated above the basal detachment (Alsop *et al.* 2017a).

Are re-fold patterns formed in slumps similar to those in metamorphic rocks?

Within orogenic belts, re-fold patterns were originally considered to be created by punctuated episodes of overprinting deformation. Each deformation 'phase' was considered to have a distinct style and orientation that were correlated over large distances within metamorphic rocks, and resulted in re-fold patterns where fold phases interfered with one another (see the discussion in Fossen 2016, p. 456; Fossen *et al.* 2018). The strict adherence to such 'D-number' schemes started to break down with the realization that heterogeneous strain could produce variable styles and orientations of structures during single progressive deformation events associated with orogenesis (e.g. Coward & Potts 1983; Holdsworth 1990; Alsop & Holdsworth 1993).

A similar range of re-fold geometries may be produced within slumps as observed in metamorphic rocks (e.g. Tobisch 1984; Farrell & Eaton 1987), although refolds within slumps form within a single slump event created by slope failure (Fig. 4a–j). Refolding that produces coaxial interference marked by Type 3 'hook' patterns may be associated with consistent directions of flow, and a late phase of pure-shear shortening that was superimposed as the slump motion reduced or stopped (Farrell & Eaton 1987). Alternatively, non-coaxial refolding marked



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by Type 1 'dome and basin' or Type 2 'mushroom' refolds implies either that the flow direction is not constant within the slump, or differential shear has locally developed within a broadly uniform slump direction leading to the creation and re-working of variably orientated folds within the flow perturbation (e.g. Alsop & Holdsworth 1993). There is no evidence within the case study that refolds are created during later loading of the slump by overburden, as the refolds are themselves truncated by erosive surfaces associated with deposition of the overlying sedimentary cap (Figs 4j & 10). Some reworking of slump folds linked to 'relaxation' and reorganization of the MTD may locally develop, and have been described by Alsop & Marco (2011, p. 449). There is also some evidence that refolding of slump folds developed within <10 cm of the sediment surface may be created by the movement of the overlying water column as it 'sloshes' back and forth across the narrow basin during seiche events (Alsop & Marco 2012b). However, overall energy within the gravity-driven MTD will generally dissipate after the initial seismic trigger and the resulting downslope movement. In summary, the observation that a similar range of re-fold styles may be created by slump folds during SSD, compared to metamorphic folds formed during HRD, demonstrates that progressive deformation is a valid mechanism to generate such interference patterns irrespective of setting.

What range of fold styles are created during slumping?

This study has shown that all potential fold styles ranging from Class 1A, 1B and 1C, Class 2, and Class 3 may form within slump systems (Figs 5a, i & 6a–h). The development of more parallel folds (Class 1B) in competent layers, and Class 3 folds in adjacent weaker layers (e.g. Fig. 6g, h) may combine to create overall similar fold packages that 'can extend for a significant distance in the profile section, in a direction parallel to the axial plane' (Price & Cosgrove 1990, p. 316). Different alternating fold styles in the bilaminate (e.g. Fig. 6g, h) may therefore encourage harmonic folding. Ortner & Kilian

(2016) also recognized fold styles that vary between Class 1B and Class 3 in adjacent layers within un lithified carbonates undergoing downslope-directed creep. Ramsay (1974) realized that if competent layers were thicker than the average in the multilayer, then accommodation structures must develop around this thicker layer, with adjacent layers displaying 'keel-like' hinges (Price & Cosgrove 1990, p. 320). Similar structures are observed in fold hinges formed next to thicker detrital layers in the present study (e.g. Fig. 8p).

Waldron & Gagnon (2011) suggested that the style and class of folds in different lithological layers may be critical in identifying SSD. Unlithified muds may display buckle folds (Class 1B), while adjacent sands are weaker (due to greater porosity and water content) and may display Class 2 (similar) folds. These relationships in unlithified sediments are the opposite to those typically encountered in lithified rocks, where sandstones are normally more competent than mudstones during folding associated with HRD. Indeed, Hibbard & Karig (1987, p. 848) recorded Class 1B folds in sandstones and Class 2 or Class 3 folds in mudstones, and suggested that beds may have been lithified at the time of folding in a variably deformed accretionary complex in SW Japan. In summary, although folding in unlithified sediments is formed via deformation associated with 'hydroplastic' particulate flow, rather than recrystallization as in metamorphic hard-rock deformation (HRD), the range of geometries produced are similar to one another. This includes the development of axial-planar fabrics in slump folds as observed previously within the Lisan Formation (e.g. Alsop & Marco 2014; Alsop *et al.* in press *a, b*). Different lithologies may, however, behave more or less competently during HRD or SSD and thereby create folds of distinct styles in each setting.

How are slump folds subsequently modified?

Post-buckle flattening. Ramsay (1967, p. 411) suggested that Class 1B parallel folds may be subsequently modified by 'flattening' that resulted in Class 1C geometries. Ramsay (1967, p. 433) went

Fig. 7. Analysis of post-buckle flattening in folds developed directly beneath sedimentary caps. In each case, paired photographs and associated inverse layer thickness analysis based on Lisle (1992) are shown in (a–b) from Peratzim, (e–f) from Masada and (i–j) from Peratzim. (a) & (b) illustrate the inverse thickness method, which plots $(1/t)$ for various orientations of the layer tangent around the fold, where t is the orthogonal thickness between tangents for the inner and outer layer boundaries (see Lisle 1992, p. 370). The inverse thickness $(1/t)$ is plotted from a common central point, each in the direction of the tangent to create an array of points to which a best-fit ellipse is matched. Examples of upslope (SW in red) tangents drawn at 70° and downslope (NE in blue) tangents drawn at 45° are shown in (a) & (b). Photographs and ellipses are 'unstrained' back to circles using the method of Srivastava & Shah (2006), and are shown in the right-hand side in (c) & (d), (g) & (h) and (k) & (l). In each case, data collected from the western or SW (upslope) limb are shown in red, while those data collected from the eastern or NE (downslope) limb are shown in blue. The change in area from ellipse to restored circle does not represent a change in volume of the sediment.

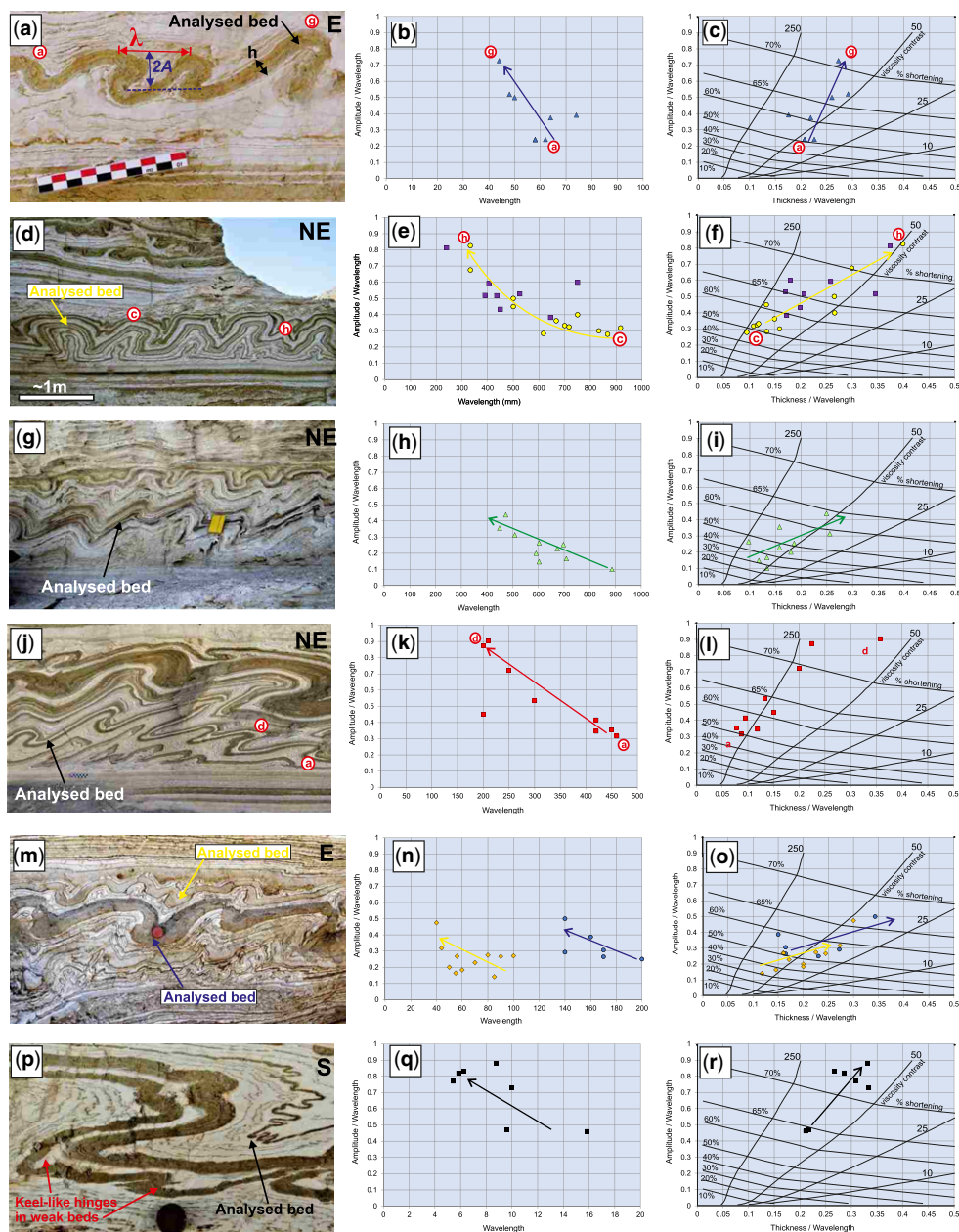


Fig. 8. Photographs and associated graphs showing wavelength plotted against amplitude/wavelength, and strain contour maps of *Schmalholz & Podladchikov (2001)* for a series of individual fold trains (a)–(c), (d)–(f), (g)–(i), (j)–(l), (m)–(o) and (p)–(r). In (e) & (f) blue squares represent data collected from the same bed in an adjacent outcrop of the slump. In (a) the wavelength (λ) is defined as the distance between two points that occupy a similar position on the fold train (i.e. between adjacent synform hinges). The thickness of a layer (h) is measured orthogonal to the folded bed, while amplitude (A) is defined as half the distance from the trough to the crest of upright folds. Strain contour plots of *Schmalholz & Podladchikov (2001)* show the estimated percentage of shortening and viscosity contrasts for folded layers. In each of the graphs, different coloured symbols represent different individual fold trains. Arrows show the general trends of data on each graph, while circled letters correspond to individual folds marked on adjacent photographs. The yellow notebook (21 cm long), the 10 cm-long chequered rule and the 15 mm-diameter coin act as scales on the photographs.

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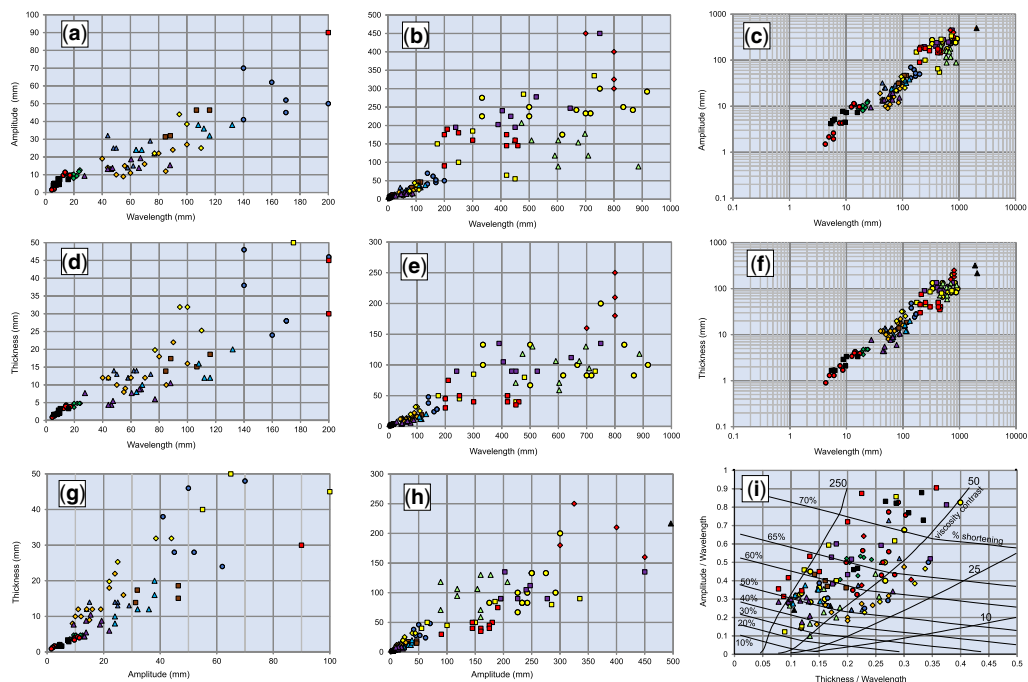


Fig. 9. (a)–(c) Graphs showing fold amplitude plotted against fold wavelength for: (a) minor folds (<200 mm wavelength), (b) mesoscopic folds (<1000 mm wavelength) and (c) overall data on a log–log plot. (d)–(f) Graphs showing folded layer thickness against fold wavelength for: (d) minor folds (<200 mm wavelength), (e) mesoscopic folds (<1000 mm wavelength) and (f) overall data on a log–log plot. (g)–(h) Graphs showing folded layer thickness against fold amplitude for: (g) minor folds (<100 mm amplitude) and (h) mesoscopic folds (<500 mm amplitude). (i) Strain contour map of *Schmalholz & Podladchikov (2001)* showing estimated percentage shortening and viscosity contrasts for folded layers. In each of the graphs (a)–(i), different coloured symbols represent different individual fold trains, including those shown in *Figure 8*.

on to suggest that ‘Many of the apparently similar folds seen in naturally deformed rocks probably owe their initial development to the buckling mechanism’. We have therefore undertaken an analysis of such post-buckle modification in order to test if this is an appropriate mechanism to create the range of fold styles within the case study. Homogeneous flattening normal to the axial plane of an original Class 1B fold results in an increased fold amplitude and a Class 1C fold (e.g. *Farrell & Eaton 1987*, p. 193), whereas flattening parallel to the axial plane leads to a reduction in fold amplitude and a Class 1A fold (e.g. *Twiss & Moores 2007*, p. 374). *Hudleston (1973)* recognized that Class 1C folds may be produced during simultaneous buckling and flattening, and there is therefore no necessity for separate and sequential ‘phases’ of fold development, with folds created during progressive deformation within the slump.

Our analysis reveals that if flattening were the sole cause of the different fold styles, then it would need to form a significant element in the

deformation. Using the technique of *Lisle (1992)* described previously, elliptical ratios of *c.* 1.7 are calculated, suggesting a relative vertical shortening of 30% (*Fig. 7a–d*). In other cases, some Class 1C folds defined by detrital-rich markers would in fact suggest up to *c.* 50% homogenous shortening (*Fig. 7i–l*). The ‘mismatch’ of axial surfaces that ‘jump’ when traced across weak beds may be a consequence of such flattening (e.g. *Fig. 3a*). However, recumbent crenulations observed in aragonite laminae in the hinges of upright Class 1A folds record only limited vertical shortening (see *Alsop & Marco 2011*). Such horizontal crenulations were attributed to ‘compaction’ by *Alsop & Marco (2011)*. As there is no other evidence for significant flattening, with thrusts preserving pristine dip angles of 30° (e.g. *Alsop et al. 2017a, b*), and a complete absence of folded clastic dykes (*Alsop et al. in press a*), this amount of vertical shortening seems unfeasible. Post-buckle flattening may represent a small component that affects upright Class 1B folds close to the sediment surface where uncompacted sediment was most water-saturated,

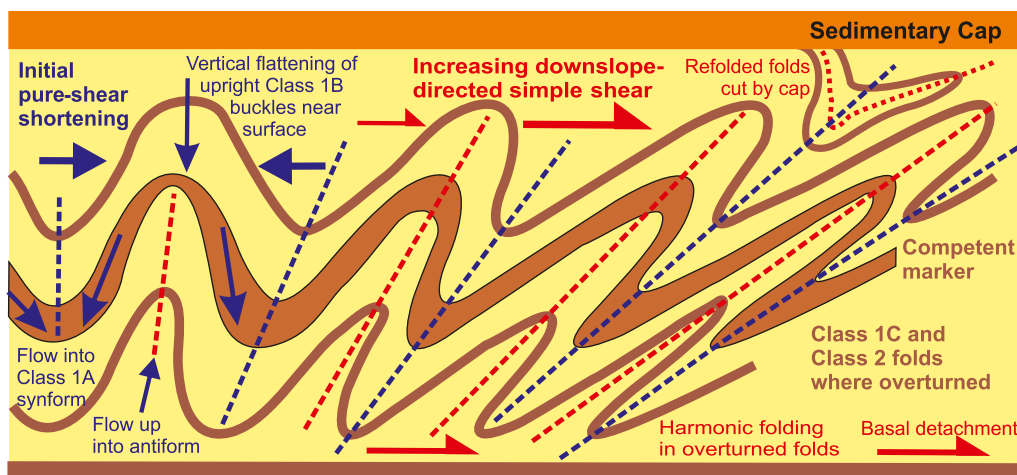


Fig. 10. Schematic summary cartoon highlighting the main geometries and styles of folding observed during soft-sediment deformation of competent (brown) markers above a basal detachment. Layer-parallel shortening associated with pure shear creates upright Class 1B buckles (left of diagram) that undergo vertical flattening (large blue arrow) and flow of sediment down fold limbs towards synformal troughs to create Class 1A antiforms. A potential component of density-driven flow up into antiforms may be associated with Rayleigh–Taylor instabilities. Increasing downslope-directed simple-shear deformation towards the right of the figure results in rotation of antiformal and synformal axial surfaces (red and blue dashed lines, respectively) towards the flow plane, resulting in overturned and harmonic folds with Class 1C or Class 2 (similar) geometry. Refolded folds created during the progressive downslope-directed shearing are truncated by the overlying erosive sedimentary cap (orange bed) that is deposited from suspension following slope failure.

and results in locally thickened limbs and thinned hinges (Class 1A) (Fig. 10). However, given the lack of supporting evidence, post-buckle flattening should not perhaps be applied to the entire sequence and we therefore suggest that other mechanisms may also have operated to modify fold shapes.

Post-buckle shearing. The relationship between layer thickness and fold wavelength suggests buckling is the dominant fold mechanism within slump folds of this case study and elsewhere (e.g. Woodcock 1976b; Farrell & Eaton 1987). As noted above, deviation of buckle folds from Class 1B may be a consequence of vertical flattening, creating upright Class 1A with thickened limbs (e.g. Figs 6a, b & 10), or recumbent Class 1C folds with thinned limbs (e.g. Fig. 6c, d) (e.g. Farrell & Eaton 1987, fig. 11). However, the variation in geometry may also be created by simple shear being imposed on the buckles, resulting in thickened or thinned fold limbs depending on the orientation of the fold and amount of shear (e.g. Farrell & Eaton 1987; see the discussion in Twiss & Moores 2007).

It has previously been suggested that buckle folds within the Lisan Formation may form during layer-parallel shortening associated with pure shear, and that fold initiation may also potentially involve density-driven flow into antiforms associated with

Rayleigh–Taylor instabilities (e.g. Alsop & Marco 2011, p. 453) (Fig. 10). Downslope-directed simple shear is then progressively imposed on these folds to create the observed range of geometries (Alsop & Marco 2011, p. 446) (Fig. 10). The tightening and rotation of upright buckles to overturned folds in the case study (Fig. 5e, f, g) is consistent with increasing downslope-directed non-coaxial-dominated deformation (e.g. Woodcock 1976b; Farrell & Eaton 1987; Twiss & Moores 2007, p. 360; Dasgupta 2008) (Fig. 10). This interpretation is also supported by numerical simulations of downslope-directed gravity sliding that results in buckle folds with hinges that rotate towards the flow direction and axial planes that sequentially rotate into the plane of flow during continued slumping (Schmalholz & Schmid 2012, p. 1817).

Recumbent folds are sometimes marked by thicker overturned lower limbs compared to the upper fold limbs (e.g. Figs 3a & 5h, i). Folds with thicker lower limbs suggest that the original buckle was upright, and that the downslope limb has spent some time in the contractional field of the strain ellipsoid (while the upslope limb was always in the extensional field and therefore thinner). In other cases, the lower fold limbs are more thinned (Fig. 6g, h), suggesting that these buckles may have originated with a downslope asymmetry with shorter,

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steeper limbs marking vergence towards the basin (see [Alsop & Marco 2013](#), fig. 7). The observation that most folds in the case study display thinned lower limbs is consistent with original buckles displaying a downslope-directed asymmetry and vergence towards the basin. In addition, we do not generally observe recumbent Class 1B or upright Class 1C or Class 2 folds, suggesting that folds are modified as they become overturned during downslope-directed shearing. These observations collectively suggest that slump folds originated as more upright buckles that have been subsequently deformed by downslope-directed simple-shear deformation (see [Farrell & Eaton 1987](#); [Alsop & Marco 2011](#)) (Fig. 10). There may therefore be a combination of post-buckle flattening and shearing, with recumbent Class 1C being created by flexural flow rather than by the flattening of original Class 1B buckles.

Post-buckle flow. [Price & Cosgrove \(1990, p. 402\)](#) noted that ‘post-buckle flattening may involve physical migration of the relatively incompetent material’. They further noted ([Price & Cosgrove 1990, p. 403](#)) that incompetent layers in the short limbs of folds may have thinned due to ‘the migration of material from the limbs to the hinge region in response to stress gradients, which tend to become greater during the last stages of fold amplification’. Although our analysis is performed on unlithified detrital-rich beds, these beds are exceptionally weak and capable of flow when compared to lithified rocks.

Within the case study slump folds, some upright synforms are markedly thickened in the hinges when compared to adjacent antiforms (Figs 7e & 8m), while the converse is rarely observed. This suggests that in some cases there may be a component of material flow from the antiformal crest down into the synformal troughs. The typical position of Class 1A folds towards the upper parts of slumps, directly below erosive truncation surfaces linked to the sediment–water interface, suggests that gravity-driven migration of sediment within steep water-saturated fold limbs close to the surface may be a significant factor (e.g. Figs 7 & 10). [Farrell & Eaton \(1987, p. 186\)](#) examined slump folds in Cyprus and the Spanish Pyrenees, and also noted the ‘flowage of sediment into cores of folds’, which on examination of their figure 3d is predominantly towards the troughs of synforms. Flow of sediment down off earlier structural highs and folds was previously termed ‘relaxation’ by [Alsop & Marco \(2011\)](#), and may form a significant local factor in fold modification.

Summary. A range of post-buckle processes may have contributed to the final geometry of the slump folds. Thickening of upright fold limbs to create

Class 1A folds suggests that a component of post-buckle vertical ‘flattening’ has operated during the slump process, while thickening of synformal hinges relative to neighbouring antiforms is consistent with gravity-driven hydroplastic flow of sediment from crests to troughs (Fig. 10). Although buckle folds may have initiated during layer-parallel shortening associated with pure shear, they were subsequently modified by downslope-directed simple shear, resulting in overturned Class 2 folds ([Farrell & Eaton 1987](#); [Alsop & Marco 2011, 2013](#)) (Fig. 10). In addition, other more nebulous influences such as locally increasing fluid pressure during the folding process may lead to a relative weakening of competent layers (e.g. [Price & Cosgrove 1990, p. 295](#)), thereby encouraging more Class 1C and Class 2 similar folds to develop during the evolution of the slump. Overall volume loss via the expulsion of fluids (e.g. [Price & Cosgrove 1990](#)) and increased grain packing may also influence the geometry of folds, although is difficult to quantify. These processes may collectively modify the initial fold shape to create the range of fold geometries that are currently observed in the study area.

Are estimates of viscosity contrasts in slumps folds similar to those in metamorphic rocks?

[Schmalholz & Podladchikov \(2001\)](#) demonstrated via strain contour maps that viscosity contrasts between layers of between 50 and 250 are typical in metamorphic rocks ([Schmalholz & Podladchikov 2001](#); see also [Hudleston & Treagus 2010](#)). These values, however, are observed to vary depending on lithological controls (e.g. [Druguet *et al.* 2009](#), fig. 3). [Schmalholz & Podladchikov \(2001, p. 211\)](#) also noted that the strain contour map is applicable not just to buckles created during pure shear, but also to asymmetrical folds that formed during simple shear. In general, the values of viscosity contrast calculated from the slump folds in the study area display a similar range to those obtained from metamorphic rocks (Figs 8 & 9).

Thus, similar geometries between metamorphic folds created during HRD and slump folds formed during SSD are borne out by a similar range of viscosity contrasts on strain contour maps. However, folds within the case study slumps formed before the overlying sedimentary cap was deposited out of suspension ‘in a matter of just hours or days’ ([Alsop *et al.* 2016, p. 80](#)) and absolute viscosities in slumps will therefore be much lower. As a guide, samples of modern Dead Sea sediment yield viscosity measurements of just 0.3 and 3 Pa s at depths of 10 and 30 cm, respectively, below the present sediment surface ([Wetzler *et al.* 2010, p. 304](#)), while [Levi *et al.* \(2008\)](#) suggested that detrital layers in the Lisan Formation have dynamic viscosity values in the range of

0.03–0.3 Pa s. In terms of strain rates, Price & Cosgrove (1990, p. 369) estimated the time taken during HRD to create folds of 500 m wavelength, based on typical plate tectonic velocities of $c. 3.3 \text{ cm a}^{-1}$, and suggested strain rates of 10^{-13} s^{-1} . If we make a crude assumption within the case study that a 1 m wavelength slump fold is formed in 1 day (Alsop *et al.* 2016), then suggested strain rates during SSD are of the order of 10^{-5} s^{-1} , while a similar fold forming more slowly over 1 week gives strain rates of 10^{-6} s^{-1} ; then the conclusion is that, despite the geometrical similarity and a similar range of viscosity contrasts between metamorphic folds and SSD folds, the strain rates within slumps that create MTDs are 7–8 orders of magnitude greater and may be viewed as geologically instantaneous.

There are a number of considerations when estimating viscosity contrasts and total shortening within beds defining slump folds. Price & Cosgrove (1990, p. 299) noted that the amount of bulk layer-parallel shortening that develops before buckles start to amplify increases as the viscosity contrast decreases. This idea was developed further by Treagus (1997), who suggested that there may be between 5 and 20% bulk (layer-parallel) shortening before buckles start to amplify, with this pre-buckle shortening most significant in modelled folds with lower viscosity ratios (Treagus 1997; Hudleston & Treagus 2010, p. 2059). Teixell & Koyi (2003) and Koyi *et al.* (2004) also noted that $c. 20\%$ layer-parallel shortening may develop in models undergoing folding and thrusting. The net result may be that Schmalholz & Podladchikov (2001) strain plots underestimate total shortening as they fail to take account of this layer-parallel component.

It has also been suggested that for a regular fold wave train with a dominant wavelength to develop, the viscosity ratio should be $>100:1$, otherwise layer-parallel shortening (noted above) that predates buckling becomes so large that it masks and would ‘overshadow folding’ (Biot 1957; see the discussion in Price & Cosgrove 1990, p. 281). However, our results suggest that in some cases, fold wave trains may display viscosity ratios of significantly less than $100:1$ (e.g. Fig. 8c), and in some instances less than $50:1$ (Fig. 8o). It has been recognized from experiments that with increased homogeneous layer-parallel shortening, which accompanies the buckling, the wave trains with dominant wavelengths may indeed form at significantly lower viscosity contrasts of $25:1$ (Hudleston 1973). Our results are therefore consistent with a component of homogeneous layer shortening that accompanies the buckling of competent layers during the slump process. Further details of the role of viscosity ratios in the folding process are provided by Schmid *et al.* (2010) and, more recently, are reviewed by Schmalholz & Mancktelow (2016).

The correlation we describe between fold wavelength (λ) and amplitude (A) (e.g. Fig. 9a–c) is in general agreement with Wetzler *et al.* (2010), who correlated A^2 and wavenumber (i.e. $1/\lambda$) in slump folds from the Lisan Formation. Wetzler *et al.* (2010) suggested that the observed folds form part of a continuum of structures that initiate as linear waves during Kelvin–Helmholtz instability and evolve into folds, finally culminating in breccias associated with turbulence and instability (see also Alsop & Marco 2011, p. 453). Analysis of fold wavelength (λ) to layer thickness (h) from the measured slump folds (e.g. Fig. 8a) suggests a general ratio of $c. 5:1$ (Fig. 9d–f). This is significantly less than the $27:1$ ratio calculated by Twiss & Moores (2007, p. 562) for folds in lithified rocks. Although this λ/h value is viewed as a constant, it is recognized that wavelength is not a very sensitive function, as slight changes result in significant differences in the estimate of the viscosity contrast. In addition, if there is a component of multilayer folding (i.e. layers are more closely spaced), then the wavelength produced will be less than for an equivalent thickness of a single-layer buckle fold (Price & Cosgrove 1990, p. 310). The net effect of this multilayer folding would be to increase both the A/λ and h/λ values. This will result in an overestimate of the percentage of contraction, and a potential over- or underestimate of the viscosity contrast on the Schmalholz & Podladchikov (2001) strain maps designed for single-layer folds. The role of multilayer spacing on the λ/h ratios of individual folded layers has recently been summarized by Schmalholz & Mancktelow (2016, p. 1437).

It is interesting to note that buckle fold layers that are entirely isolated within aragonite-rich units may display a greater viscosity contrast as the percentage of contraction increases (e.g. Fig. 8a–c); whereas if layers are closer together, there is an apparent reduction in viscosity contrast with increasing percentage of contraction, resulting in less steep trends on strain contour plots (Fig. 8f, i, l, o). This apparent discrepancy may reflect a number of variables. If a measured single-layer fold actually forms part of a multilayer, then the multilayer will act as a thicker mechanical unit, with a greater effective thickness, meaning that measured h/λ values would be underestimated. This multilayer influence may be most significant where folds initiate and amplitudes are lower (with a lower contraction percentage) and may result in some datasets defining less steep trends on strain plots: that is, the effect of multilayers may be to overestimate viscosity contrasts at lower values of percentage of shortening (see also Druguet *et al.* 2009, p. 503). An alternative interpretation would be that viscosity contrasts actually vary as folds develop due to the expulsion of fluids during folding. Price & Cosgrove (1990, p. 398) examined folds

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formed during HRD and argued that 'the fold acts as a pump which, if the external energy is sufficient, continues to extrude fluids from itself until it is almost dry'. Any expulsion of fluids leading to variations in viscosity contrasts would have the greatest effect on resulting fold geometries if it occurred relatively early in the folding process when amplitudes are smaller and before fold wavelengths are established in the system. We here suggest that during SSD, a preferential expulsion of fluids from weaker (aragonite-rich) layers as the percentage of shortening increases would lead to an increase in the viscosity of these layers, and consequently a reduction in the contrast between the viscosities of the aragonite- and detrital-rich layers.

In summary, the overall correlation between bed thickness (h), fold wavelength (λ) and amplitude (A) suggests that buckling is the dominant fold mechanism in the slump folds of the study area (Figs 8 & 9). Strain maps developed by Schmalholz & Podladchikov (2001) are equally applicable to folds created during HRD of metamorphic rocks and slump folds formed during SSD, and suggest similar viscosity contrasts ranging from 50 to 250. Within slump folds, components of potential layer-parallel shortening would lead to an underestimate of total shortening, while an element of multilayer folding may lead to an overestimate of the percentage of contraction and less accurate estimates of viscosity contrasts. We speculate that it may also be possible for viscosity contrasts to actually vary as fold trains evolve due to the preferential expulsion of fluids as folds tighten.

Conclusions

In this study, classical techniques of structural analysis developed by John Ramsay, amongst others, during the investigation of folds formed in metamorphic terranes have been applied to folds generated within unlithified soft sediments. Using superbly exposed slump folds created during SSD within MTDs of the Lisan Formation, we have demonstrated that recumbent fold packages define an echelon fold trains that display a harmonic style of folding. Slump folds are entirely systematic and should not be viewed as 'chaotic', although upright buckles may form disharmonic folds that cannot be traced for significant distances along their axial surfaces before truncation by overlying sedimentary caps. We also recognize a similar range of refold patterns created during SSD within slump folds when compared to HRD in metamorphic rocks. Refolds may also be truncated by the overlying sedimentary cap, indicating that they formed during a single progressive slump event and are not a product of the superposition of later 'events' (Fig. 10).

A wide variety of fold styles ranging from Class 1A, 1B and 1C, Class 2, and Class 3 geometries are developed within slumps of the case study. Within individual folds, aragonite-rich layers may display Class 2 (similar) folds, while adjacent detrital-rich beds are marked by Class 1C fold geometries. Such varying styles of folding allow individual folds to propagate further along their axial planes without encountering significant accommodation problems around the fold hinge. Upright buckle folds defined by detrital-rich marker beds display a Class 1B (parallel) style of folding, while recumbent folds display Class 2 (similar) fold styles. Upright fold hinges and axial planes may undergo rotation during progressive simple-shear-dominated deformation as the MTD moves downslope under the influence of gravity. In detail, the thickening of upright fold limbs relative to the hinge to form Class 1C folds suggests a component of vertical flattening that post-dates the original folding, but predates the deposition of the overlying sedimentary cap shortly after slope failure. This flattening appears to be most significant towards the top of each MTD, where sediments are more water-saturated (Fig. 10). In addition, upright folds are also subsequently modified by the flow of sediment from antiformal crests to troughs of synforms, together with simple-shear deformation associated with continued downslope movement that results in Class 1C and Class 2 similar folds displaying recumbent attitudes (Fig. 10).

Our study has involved the measurement of fold wavelength (λ), amplitude (A) and bed thickness (h) in order to estimate the percentage of shortening and viscosity contrast between folded layers on strain contour maps (Schmalholz & Podladchikov 2001). Estimates of viscosity contrast from folded layers within metamorphic rocks formed during HRD and slump folds created during SSD both suggest broadly similar values in the range of 50–250. We speculate that it may also be possible for viscosity contrasts to vary as fold trains evolve due to preferential expulsion of fluids as folds tighten. The geologically instantaneous nature of slope failure, which may be 7–8 orders of magnitude greater than folding associated with HRD, indicates that the absolute viscosities will, however, be considerably lower in slumps.

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