Dipole in Shell

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1 Problem

What is the (vector) electric field strength at the center of a spherical cavity of radius a in a grounded conductor if a point dipole **p** is placed at distance b (0 < b < a) from the center so as to make angle α to a radius vector?



2 Solution

This problem is from Yakov Kantor's Physics Quiz site, http://star.tau.ac.il/QUIZ/

We solve via the image method for spherical geometry. Recall that a charge q at radius b < a results in an image charge q' at a^2/b , such that the potential at r = a vanishes:

$$\phi(r=a) = 0 = \frac{q}{a-b} + \frac{q'}{a^2/b-a} = \frac{q}{a-b} + \frac{q'b/a}{a-b},$$
(1)

and hence q' = -qa/b.

For the present case of a point dipole **p**, we can think of this as due to charges $\pm q$ separated by distance $\delta = q/p$, with the charge -q at radius b and the charge +q at radius $b + \delta \cos \alpha$ as shown below.

The image transformation $r \to a^2/r$ is conformal (angle preserving), so the triangle formed by the center of the sphere and charges $\pm q$ is similar to that formed by the center and the image charges q' and q''. Image charge q' = +qa/b is at distance a^2/b from the center. Hence, the distance between the two image charges is

$$\delta' = \delta \frac{a^2}{b^2} \,. \tag{2}$$



The image charge q'' has value

$$q'' = -q\frac{a}{b+\delta\cos\alpha} \approx -q\frac{a}{b} + q\delta\cos\alpha\frac{a}{b^2}.$$
(3)

Thus, taking the limit $\delta \to 0$, $q \to \infty$ such that $q\delta = p$, the two image charges can be thought of as a dipole of strength

$$p' = -q\frac{a}{b}\delta\frac{a^2}{b^2} = -p\frac{a^3}{b^3}\,,\tag{4}$$

at angle α to the radius vector as shown above, **plus** an additional charge of magnitude

$$q''' = p\frac{a}{b^2}\cos\alpha,\tag{5}$$

located at distance a^2/b from the center. If $\alpha = \pm \pi/2$, the image of the dipole **p** is a pure dipole.

To calculate the electric field at the center, we recall that the field of a point dipole \mathbf{p} is

$$\mathbf{E} = \frac{3(\mathbf{p} \cdot \hat{\mathbf{r}})\hat{\mathbf{r}} - \mathbf{p}}{r^3},\tag{6}$$

where \mathbf{r} points from the dipole to the observer. In the present problem, we define $\hat{\mathbf{z}}$ to point from the center to the dipoles, so $\hat{\mathbf{r}} = -\hat{\mathbf{z}}$. Dipole \mathbf{p} is taken to lie in the *x*-*z* plane, so that

$$\mathbf{p} = p\hat{\mathbf{z}}\cos\alpha + p\hat{\mathbf{x}}\sin\alpha,\tag{7}$$

and

$$\mathbf{p}' = p \frac{a^3}{b^3} \hat{\mathbf{z}} \cos \alpha - p \frac{a^3}{b^3} \hat{\mathbf{x}} \sin \alpha.$$
(8)

Then, the field at the center of the cavity due to the dipole, and its image dipole + extra charge q''' is

$$\mathbf{E}(0) = \frac{3(\mathbf{p} \cdot \hat{\mathbf{z}})\hat{\mathbf{z}} - \mathbf{p}}{b^3} + \frac{3(\mathbf{p}' \cdot \hat{\mathbf{z}})\hat{\mathbf{z}} - \mathbf{p}'}{(a^2/b)^3} - \frac{q'''\hat{\mathbf{z}}}{(a^2/b)^2}$$
$$= \frac{2p\hat{\mathbf{z}}\cos\alpha - p\hat{\mathbf{x}}\sin\alpha}{b^3} + \frac{2p\hat{\mathbf{z}}\cos\alpha + p\hat{\mathbf{x}}\sin\alpha}{a^3} - \frac{p\hat{\mathbf{z}}\cos\alpha}{a^3}$$
$$= p\hat{\mathbf{z}}\cos\alpha\left(\frac{2}{b^3} + \frac{1}{a^3}\right) - p\hat{\mathbf{x}}\sin\alpha\left(\frac{1}{b^3} - \frac{1}{a^3}\right).$$
(9)